



The Effects of Layer Constraint on Stress Wave Propagation in Multilayer Composite Materials

by Alper Tasdemirci, Ian W. Hall, Bazle A. Gama, and Mustafa Guden

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**Department of Mechanical Engineering
University of Delaware
Center for Composite Materials
Newark, DE 19716**

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14. ABSTRACT Multilayer materials consisting of ceramic and glass/epoxy with a rubber interlayer have been subjected to a high strain rate compression using a split-Hopkinson Pressure Bar (SHPB). The feasibility of modeling stress wave propagation in complex multilayer materials has been demonstrated. It has been shown that the effects of lateral confinement of a normally low-modulus interlayer material can significantly affect the response to wave propagation. Numerical modeling clearly shows that severe stress inhomogeneities and discontinuities exist, and these may have serious consequences for the mechanical and other properties. The one-dimensional stress state usually assumed for conventional SHPB testing is therefore inapplicable, and both numerical and experimental results have to be coupled for a complete understanding of the wave propagation characteristics. In this study, both methods were used, and the stress states inside the components were presented.					
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Contents

List of Figures	iv
List of Tables	iv
Acknowledgments	v
1. Introduction	1
2. Experiments and Modeling	2
3. Results	5
3.1 Impact Velocity (10 m/s).....	5
3.2 Impact Velocity (16 m/s).....	6
3.3 Impact Velocity (20.5 m/s).....	9
4. Discussion	13
5. Conclusions	17
6. References	18
Distribution List	21

List of Figures

Figure 1. (a) Constrained sample prior to testing and (b) schematic of setup.	3
Figure 2. Stress on the incident and transmitter bars during a test at 10 m/s (unconstrained-rubber): (a) experimental and (b) calculated.....	6
Figure 3. (a) Experimental and (b) calculated stress on ceramic: (c) experimental and (d) calculated stress on composite (unconstrained-rubber, $V = 10$ m/s).....	7
Figure 4. Stress on the incident and transmitter bars during a test at 10 m/s (constrained-rubber): (a) experimental and (b) calculated.....	8
Figure 5. Stress on the specimen (constrained-rubber) tested at 10 m/s: (a) experimental and (b) calculated stress on ceramic: (c) experimental and (d) calculated stress on composite.....	9
Figure 6. Stress measured on the incident and transmitter bars during a test at 16 m/s: (a) unconstrained rubber and (b) constrained rubber.....	10
Figure 7. Experimental data from ceramic/rubber/composite tested at 16 m/s. Stress measured on (a) ceramic and (b) composite.	10
Figure 8. Calculated output from strain gages on the incident and transmitter bars during a test at 16 m/s: (a) unconstrained rubber and (b) constrained rubber.....	11
Figure 9. Calculated data from ceramic/rubber/composite ($V = 16$ m/s). Stress measured on (a) ceramic and (b) composite.....	11
Figure 10. Stress on the incident and transmitter bars during a test at 20.5 m/s on an unconstrained ceramic/rubber/composite: (a) experimental and (b) calculated.	12
Figure 11. Stress measured on ceramic: (a) experimental and (b) calculated. Stress measured on composite: (c) experimental and (d) calculated (unconstrained-rubber, $V = 20.5$ m/s).	13
Figure 12. Stress measured on the incident and transmitter bars during a test at 20.5 m/s on a constrained ceramic/rubber/composite: (a) experimental and (b) calculated.	14
Figure 13. Stress measured on ceramic: (a) experimental and (b) calculated: Stress measured on composite: (c) experimental and (d) calculated (constrained-rubber, $V = 20.5$ m/s).	15

List of Tables

Table 1. Material properties used in finite element models.....	5
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1. Introduction

Thick-section composite materials are frequently used under dynamic loading conditions, but their behavior is still not clearly understood. Impact loading of monolithic laminates has been the subject of several investigations, e.g., with glass/epoxy and graphite/epoxy (1–10). Similarly the penetration or perforation of composites has also been studied (11, 12), but severe complications arose whenever widely dissimilar materials were in intimate contact because their differing impedances caused complex wave reflection and transmission phenomena at each interface encountered. Thick, layered, or graded structures have significant potential for armor applications, and Li et al. (13) reported the dynamic characterization of layered and graded structures under impulsive loading. Another example of multilayer materials is provided by modern integral composite armor for vehicle applications as described by Fink (14) and Gama et al. (15–17). The armor material must provide ballistic protection at minimum weight and may contain several layers of different impedance, usually a ceramic layer followed by a thick composite plate (e.g., glass fiber/epoxy). High-velocity impact of this type of integral armor has been the subject of finite element studies by Mahfuz et al. (18). Jovicic et al. (19) modeled the ballistic behavior of gradient design composite armors.

The elastic adhesives used in composite armors can also alter wave propagation in the armor components. The mechanical behavior of different elastic adhesives under impact loads was studied by Martinez et al. (20) who reported that the capability of transmitting and reflecting the impact energy depends on the thickness and the type of the adhesive used. They concluded that the utilization of a thin layer of a rigid adhesive was the best way to transmit energy with the lowest reflection coefficient.

A central concept in enhancing the ability of multilayer material structures to withstand rapid impulsive loading is to spread the local impact load as rapidly and widely as possible. This can be achieved by placing a high wave-speed layer in the layered system. Gupta and Ding (21) studied numerically the effects of wave speed, layering geometry, and mechanical properties of the layer and substrate on load spreading. They showed that for a fixed layer thickness, a single thick high-strength high wave-speed layer appears to be able to offer the best lateral load spreading through intense and rapid wave transmission and spreading. The low wave-speed material used in multilayered targets appears to deteriorate the load spreading capability of the layered system.

Design of efficient multilayer materials for impact resistance requires both modeling and experimental efforts, and the split-Hopkinson pressure bar (SHPB) is a convenient tool in the latter respect although conventional data reduction routines obviously cannot be used for these materials. Two- and three-dimensional (2-D and 3-D) wave propagation in Hopkinson bar tests

has been investigated numerically by several authors (22, 23). In experimental studies, impact velocity of the striker bar and axial strain on the bar surfaces are the most commonly measured quantities. In numerical studies, besides the parameters previously mentioned, displacement and velocity of nodes, strain and stress of the elements, and interface forces can all be acquired as a function of time (24).

Prior work (25) addressed the situation with three different layers consisting of ceramic, rubber and composite, in which lateral expansion of the rubber was permitted. However, in practical large-scale structures, the rubber interlayer would be constrained by the surrounding material, and this present report considers the effects of such lateral constraint on the resulting properties. Samples used in SHPB testing can, at the most, have the same diameter as the bar, and earlier experiments showed that considerable radial flow occurred in the rubber interlayer. Larger samples, typical of many anticipated applications, would be subjected to severe lateral constraints which would, in turn, affect the through-thickness stresses reported. In fact, the real case will probably lie somewhere between the extremes of completely constrained and completely unconstrained interlayers, so the evaluation of both limiting cases is correspondingly important.

The effect of lateral constraint on ballistic performance has been investigated by many authors, and it is now well known that a compressive prestress is helpful in improving the fracture energy and impact resistance in brittle materials. Espinosa et al. (26) experimentally studied the impact resistance of ceramics confined in steel fixtures and showed the enhancement of ballistic efficiency of the confined ceramics. Martinez et al. (20) have determined the stress-strain curve of confined adhesive used in armor at high strain rates. Within the armor, elastic adhesives, comparable with the rubber interlayer in our case, were used to bond two large plates of much more rigid materials, which themselves impede subsequent lateral displacements of the adhesive.

This study, then, presents the initial results of a combined experimental and numerical investigation and serves to delineate the principal features and identify the problems to be solved in order to develop a better understanding of the effects of constraint in such multilayer materials.

2. Experiments and Modeling

Samples were prepared from multilayer plates with layers of widely different impedances. The plates consisted of three layers, namely a 13.96-mm-thick alumina ceramic, a 1.5-mm-thick layer of ethylene propylene diene monomer (EPDM) rubber, and a layer of glass/epoxy composite.

The 5×5 plain weave S-2 Glass^{*} fiber woven fabric (0.814 kg/m^2) SC15[†] epoxy (toughened resin) composite plates were 11.3-mm-thick and were produced using the vacuum-assisted resin transfer molding process. Lateral confinement of the rubber interlayer was obtained by placing a 6-mm-wide steel retaining ring around the junction of the sample as illustrated in figure 1. An interference fit was achieved between the rubber and the steel ring, and no rubber was squeezed out into the region between them while testing. Possible inertial effects and interactions between the ring and the other components of the sample were checked via tests on individual ceramic and composite samples with the ring in place. No modification to the wave propagation behavior was observed in the presence of the steel ring. The ceramic layer was always at the impacted side.

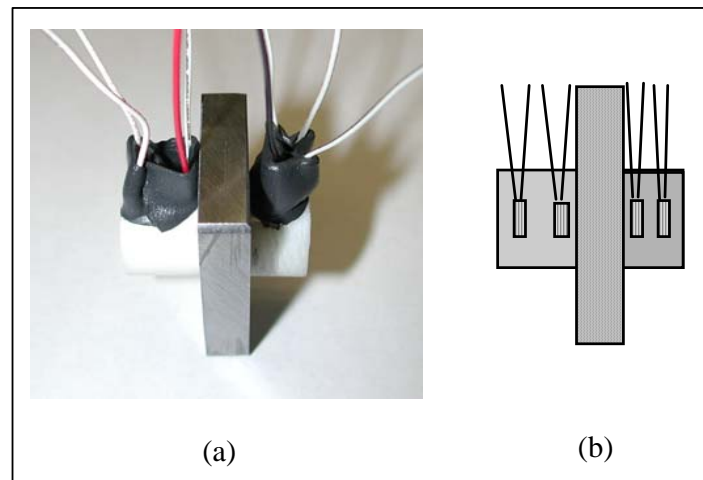


Figure 1. (a) Constrained sample prior to testing and (b) schematic of setup.

Cylindrical samples, 15.7 mm in diameter, were core drilled from the plates in the through-thickness direction. Samples were compression tested over a wide range of displacement rates using the SHPB apparatus (the compression axis normal to fiber plane). However, the focus of the present report concerns a series of tests, all of which were conducted with striker bar velocities of 10, 16, or 20.5 m/s—as an approximate guide, these velocities would generate “average strain rates” of ~ 400 , 500, and 700/s, respectively.

The particular SHPB apparatus used consists of Inconel[‡] 718 bars, a 356-mm-long striker bar, 3450-mm incident, and 1850-mm transmitter bars, all with a diameter of 19 mm. Further details of the experimental setup and standard data reduction routines are available elsewhere (27). Samples were fitted with strain gages, as shown in figure 1, so as to monitor real-time strains/stresses during the course of the tests. Strain gages with 0.79-mm element lengths were

^{*}S-2 Glass is a registered trademark of Owens Corning.

[†]SC15 is a trademark of Applied Poleramic Incorporation..

[‡]Inconel is a registered trademark of the INCO family of companies.

used generally, although several tests were also carried out with an array of gages designed to sample the strain simultaneously at several locations along the sample length and thus provide a strain/time/position map of the wave passage.

A 3-D SHPB finite element model was used to study stress wave propagation in the multilayer materials and also in the individual components. Rubber is a highly nonlinear elastic material and the role of this nonlinear material has been studied by modeling the rubber layer with experimentally determined material data. The analyses were performed using a commercial explicit finite element code LS-DYNA 960. Two axes of symmetry were assumed, so only one quarter of the bar was modeled. For each test modeled, the output was displayed at several locations, within the sample as well as at the location of the strain gages on the incident and transmitter bars of the SHPB apparatus. The desired ideal result is that output from the strain gages on the incident and transmitter bars closely match data calculated from the model. Similarly, output measured by gages on the sample should also closely match data calculated from the model. When both these conditions are met, it indicates that the model is accurately capturing the wave propagation behavior in the sample and bars.

The model has four components in contact: a striker bar of 356 mm in length, an incident bar and a transmission bar each of 1524 mm in length, and the specimen, i.e., the ceramic, rubber, and composite layers, the thicknesses of which are 14, 1.5, and 10.6 mm, respectively. The bar diameter is 19.05 mm, and the diameter of the specimen is 16.0 mm. The component materials are modeled with 8-node solid elements, and the interfaces are modeled with the automatic contact sliding interfaces without friction. The impact velocity of the striker bar ($V = 10, 16, \text{ and } 20.5 \text{ m/s}$) has been defined as the initial condition, and all other boundaries are traction free and can move in any direction. In order to save computation time, the simulation uses bars 1524 mm in length, instead of full length bars. It will be seen later from the figures that this has the effect of decreasing the transit time between successive waves and shortening the wave duration slightly, however, it does not affect the basic wave shapes or amplitudes. A few trial computations were carried out using full-length bars, but apart from the slightly smaller time window, no significant differences were found, and the shorter bars were used henceforth.

Material properties used in the finite element code are shown in table 1. The ceramic was modeled with an isotropic elastic material model, and the composite was modeled with an orthotropic elastic material. Rubber was modeled with two different material models. The Mooney-Rivlin Model (28) (two parameter nonlinear material model) was used for the unconstrained configuration and the Blatz-Ko Material Model (28) was used for the constrained configuration. The Blatz-Ko Material Model shows better agreement for the hydrostatic state of stress of the rubber. For the unconstrained case, the rubber interlayer deforms very extensively, and this large deformation caused stability problems in the finite element model. This problem was solved by using different material parameters for the different cases. While higher shear modulus values give better results for the constrained case, in unconstrained samples, the lower shear modulus values give better agreement with the experimental results. Actual compression

tests on the rubber itself confirm that this behavior is indeed observed in practice. To be able to use the same material model would probably be preferable, and this will be implemented for future simulations. The Inconel bars were modeled with an isotropic elastic material model, and lateral confinement of the rubber interlayer was modeled by preventing the displacements in both x and y directions for this component.

Table 1. Material properties used in finite element models.

Material	Modulus of Elasticity (GPa)	Poisson's Ratio	Density (kg/m ³)	Other
Ceramic	370	0.22	3900	—
Mooney-Rivlin rubber	—	0.495	1200	A: 0.2 (MPa) B: 0.8 (MPa)
Blatz-Ko rubber	—	0.493	1200	G: 20 (MPa)
Composite	E ₁ : 40 E ₂ : 40 E ₃ : 15	v ₂₁ : 0.12 v ₃₁ : 0.173 v ₃₂ : 0.173	1668	G ₁ : 8 (GPa) G ₂ : 8 (GPa) G ₃ : 8 (GPa)
Inconel	207	0.3	7850	—

3. Results

The experimental results are presented here in order of increasing incident bar velocity which corresponds to increasing loading rate and, as will become clear, increasing degrees of damage within the samples. Three different striker bar velocities were used, and SHPB tests and simulations were performed for both the unconstrained and the constrained situations. The primary data for each test consist of (1) experimental output from the SHPB bars for constrained and unconstrained specimens, (2) measured strain gage data from each sample, and (3) numerical data, which are then compared with the corresponding experiments.

3.1 Impact Velocity (10 m/s)

Figure 2 shows experimental and calculated SHPB data from an unconstrained sample tested at a striker bar velocity of 10 m/s, and close agreement is noted between the experimental and numerical results. Experimentally, it is seen that the transmitted wave amplitude slowly increases to ~60 MPa as indicated and exhibits a minor peak of ~100 μ s before that. Calculated data show almost identical behavior.

Experimental data from the ceramic portion of an unconstrained strain-gaged sample are shown in figure 3a, while figure 3b shows the corresponding numerical data. The insets in the figure indicate the location of the gages or nodes interrogated. First, it is noted that the stress varies greatly with time and, second, it is noted that the stress close to the incident bar/ceramic interface

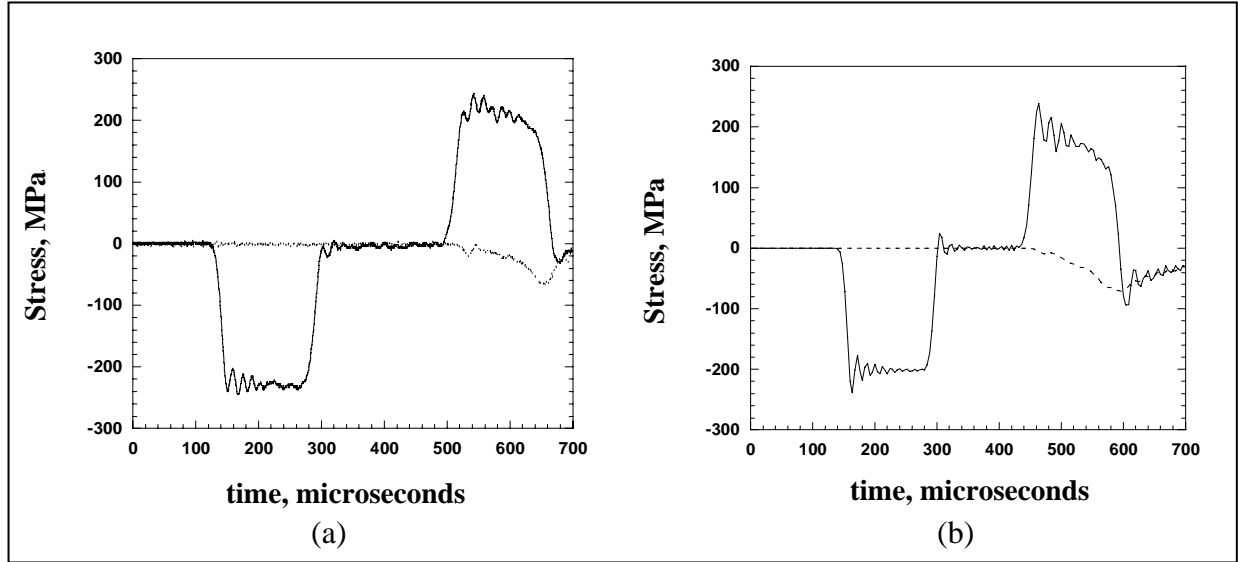


Figure 2. Stress on the incident and transmitter bars during a test at 10 m/s (unconstrained-rubber): (a) experimental and (b) calculated.

is invariably greater than that closer to the ceramic/rubber interface. Similarly, figure 3c and d shows experimental and calculated stresses measured on the composite portion of the sample.

Note that the magnitude of the initial peak in the composite layer is less than that in the ceramic layer and, again, an inhomogeneous stress distribution exists. Also, fewer large stress oscillations are noted than in the case of the ceramic. Generally, the numerical data show broadly similar behavior to the experimental data in each case, including multiple peaks in the ceramic, similar scale of stress inhomogeneity, similar magnitudes of the maximum stress, and a similar overall shape to each stress vs. time curve.

Figure 4 shows experimental and numerical SHPB data from a sample, tested at a striker bar velocity of 10 m/s, in which the rubber interlayer was constrained. Comparison with corresponding unconstrained data (figure 2a and b) shows that constraint greatly modifies the reflected and transmitted wave shapes. The constrained sample exhibits a maximum transmitted wave amplitude of ~ 200 MPa as compared to ~ 60 MPa for the unconstrained case. Figure 5 shows experimental and calculated stresses within the ceramic and composite, as a function of time, at different locations within the sample. It shows significantly different behavior compared to the unconstrained data, namely, what is essentially a single peak in the ceramic and composite and a more rapidly rising stress in both of the components when the rubber interlayer is constrained. The peak stress values in each component are almost the same value, i.e., ~ 250 MPa—considerably higher than when the rubber is unconstrained.

3.2 Impact Velocity (16 m/s)

When tested at an intermediate velocity, samples began to suffer limited damage, although none failed catastrophically. Lateral constraint of the rubber interlayer was found to increase the

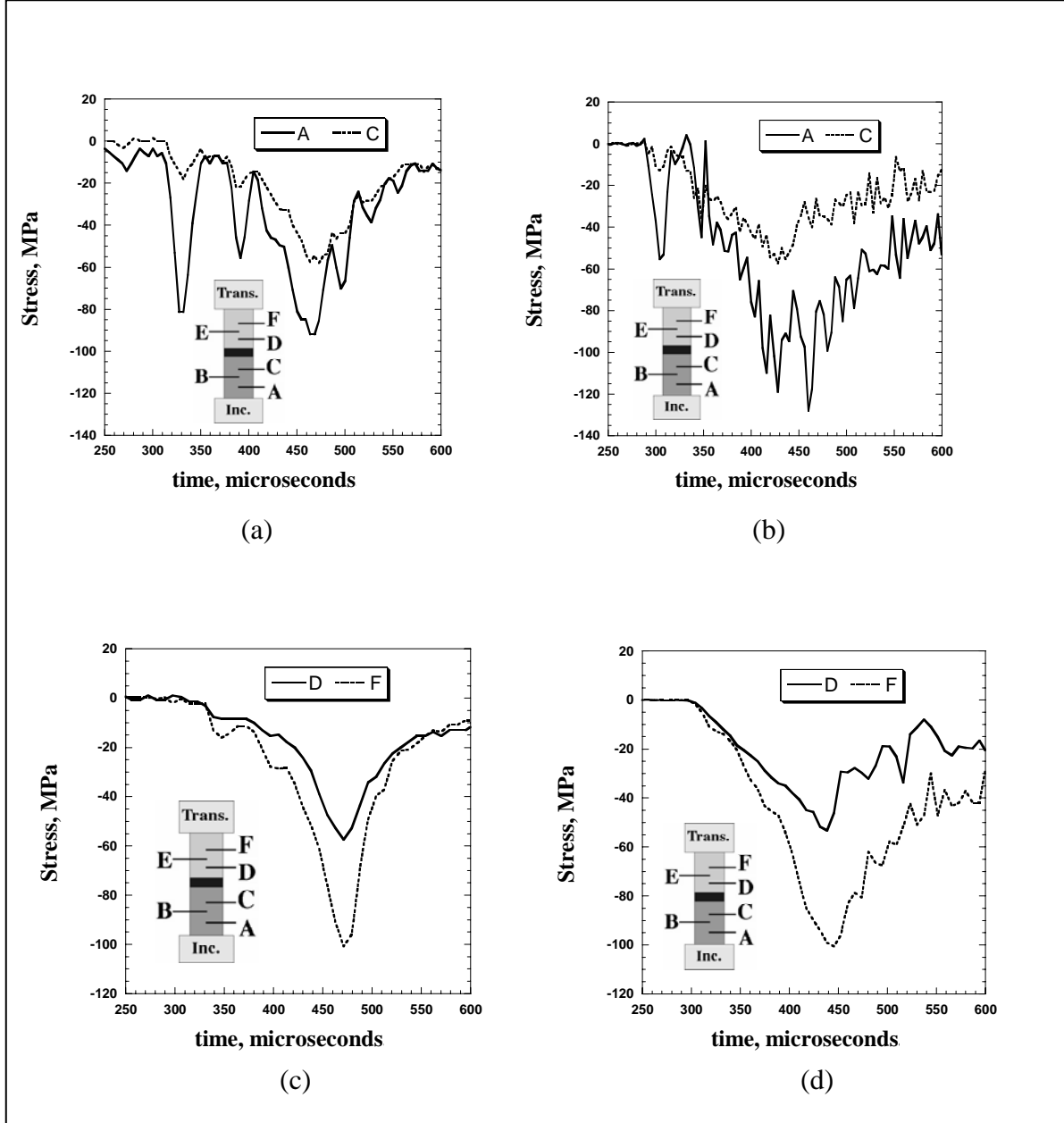


Figure 3. (a) Experimental and (b) calculated stress on ceramic: (c) experimental and (d) calculated stress on composite (unconstrained-rubber, $V = 10$ m/s).

damage level in both of the components. Visual damage in the composite exhibited itself as lateral spreading of the layers accompanied by significant radial strain, whereas the ceramic only exhibited occasional and limited spalling from the edges of the impacted face.

Figure 6a and b shows experimental SHPB data for the trilayer ceramic/rubber/composite with the rubber layer unconstrained and constrained, respectively. It is clear that the shapes of the transmitted and reflected waves for the constrained configuration are drastically different from the unconstrained case.

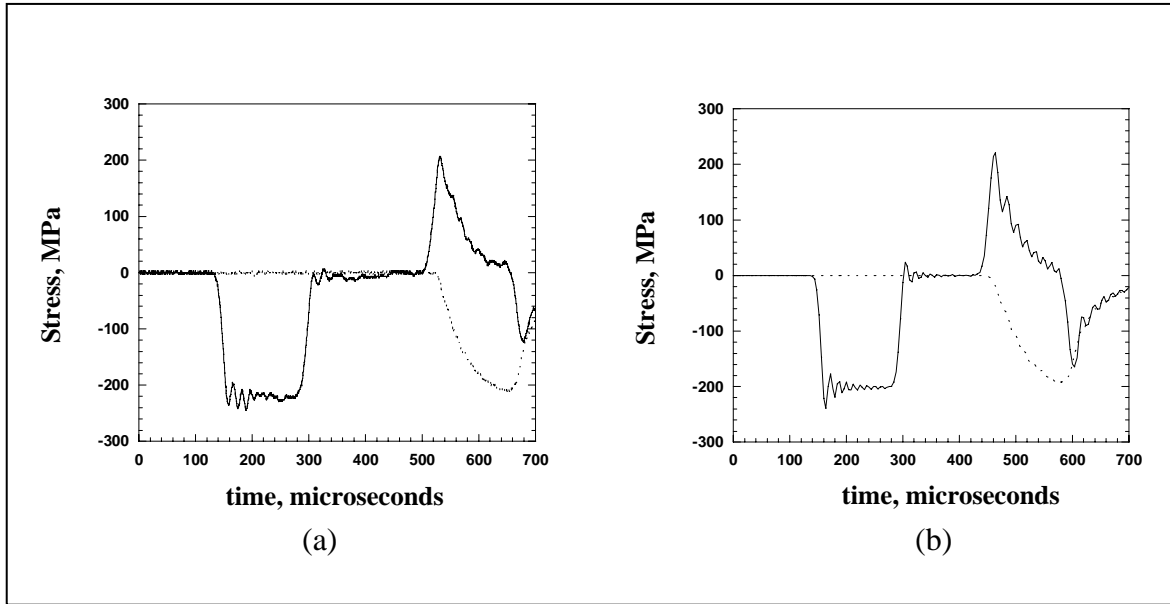


Figure 4. Stress on the incident and transmitter bars during a test at 10 m/s (constrained-rubber): (a) experimental and (b) calculated.

Major differences in the wave propagation characteristics, caused by lateral confinement of the rubber interlayer, are also clearly demonstrated in strain gage data collected from the separate layers. Figure 7a and b shows experimental data from the ceramic and composite layers, respectively. Data from gages on the ceramic show a marked increase in the measured stress levels resulting from confinement (figure 7a). Nevertheless, the general complexity of the wave forms of the unconstrained sample remains reminiscent of those tested at lower velocity (see figure 3a) insofar as three peaks may be discerned at $\sim 65\text{-}\mu\text{s}$ intervals. The peaks present in the unconstrained case merge into essentially one peak when the rubber is constrained. In the composite, the maximum stress level experienced with constraint of the rubber interlayer is $\sim 2.5\times$ that for the unconstrained sample (figure 7b), while the corresponding stress in the ceramic increases by a factor of ~ 3.3 .

Figure 8a and b shows the calculated data from the Hopkinson bars. For these samples, agreement between the experimental (figure 6a and b) and numerical data is currently slightly less close than for the low-velocity case, principally because damage has begun to occur in the experimental samples, but damage mechanisms have not yet been included in the present model, although they have been discussed elsewhere (29). The effect of damage initiation is reflected in a truncation of the early peak in the measured reflected wave.

Figure 9a and b shows numerical data from the individual ceramic and composite layers for the constrained and unconstrained cases. The elements chosen for the numerical data were at approximately the same position as the strain gages reported in figure 7. Comparison of these two figures shows that the calculated stress magnitudes are very similar to the experimental

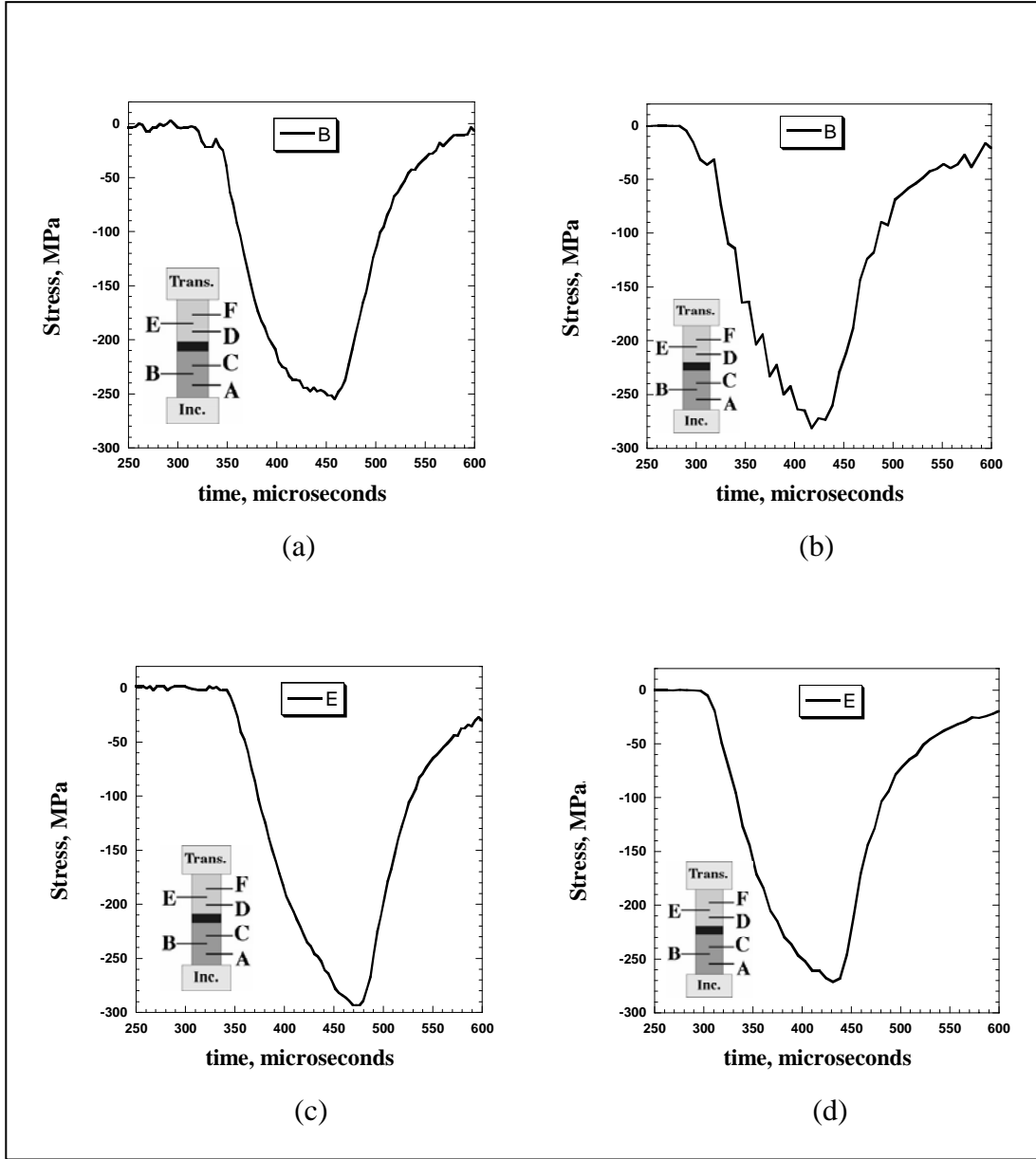


Figure 5. Stress on the specimen (constrained-rubber) tested at 10 m/s: (a) experimental and (b) calculated stress on ceramic: (c) experimental and (d) calculated stress on composite.

values. They differ only very slightly in detail as a result of the slight “stress-averaging” effect due to the finite size of the gages and, despite this, the same general shapes are found.

3.3 Impact Velocity (20.5 m/s)

A similar set of experiments and simulations was then carried out for a higher striker bar velocity. Now the presence of the rubber interlayer and its constraint leads to major differences in the wave propagation characteristics. Figure 10a and b shows experimental and calculated data from the incident and transmitter bars for the unconstrained rubber case. For this configuration, the basic shapes and magnitudes of the transmitted waves resemble each other, but

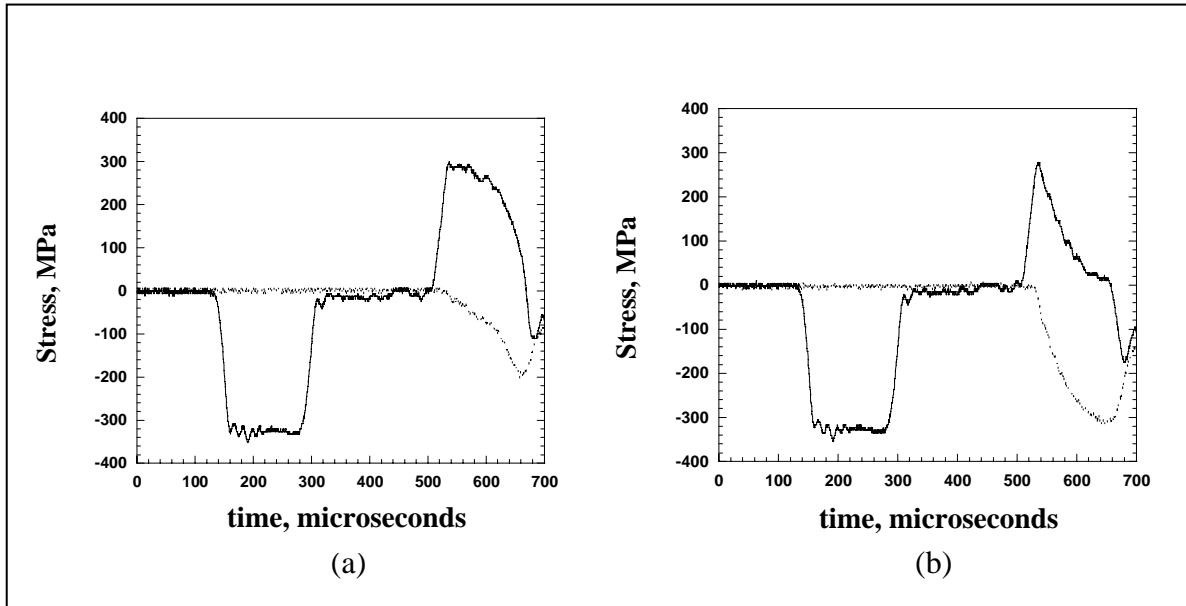


Figure 6. Stress measured on the incident and transmitter bars during a test at 16 m/s: (a) unconstrained rubber and (b) constrained rubber.

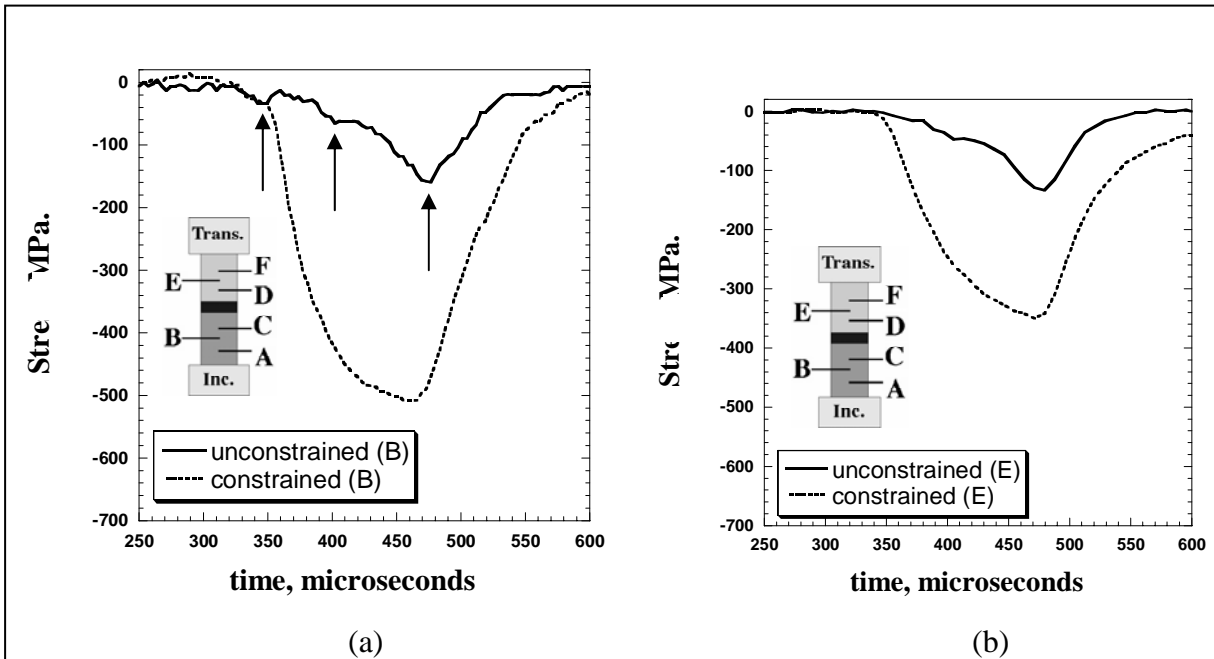


Figure 7. Experimental data from ceramic/rubber/composite tested at 16 m/s. Stress measured on (a) ceramic and (b) composite.

it is evident that significant physical damage begins to occur during the test at a stress of ~ 400 MPa, most clearly indicated by a change in the reflected wave shape.

Figure 11 shows experimental and numerical data from the individual ceramic and composite layers, respectively, with an unconstrained rubber interlayer. Figure 11a and c shows the actual

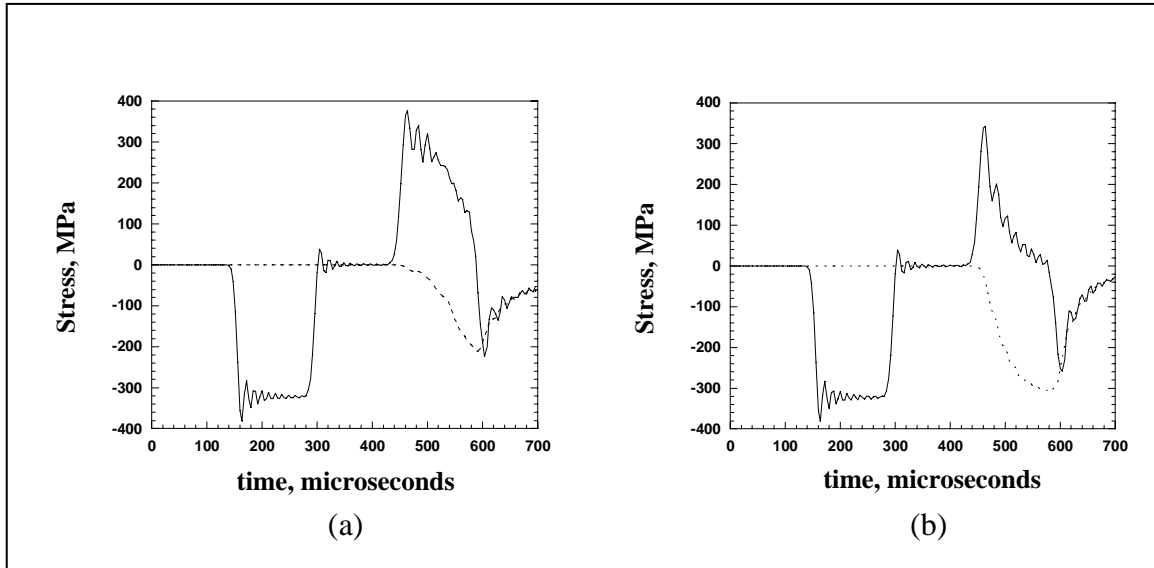


Figure 8. Calculated output from strain gages on the incident and transmitter bars during a test at 16 m/s: (a) unconstrained rubber and (b) constrained rubber.

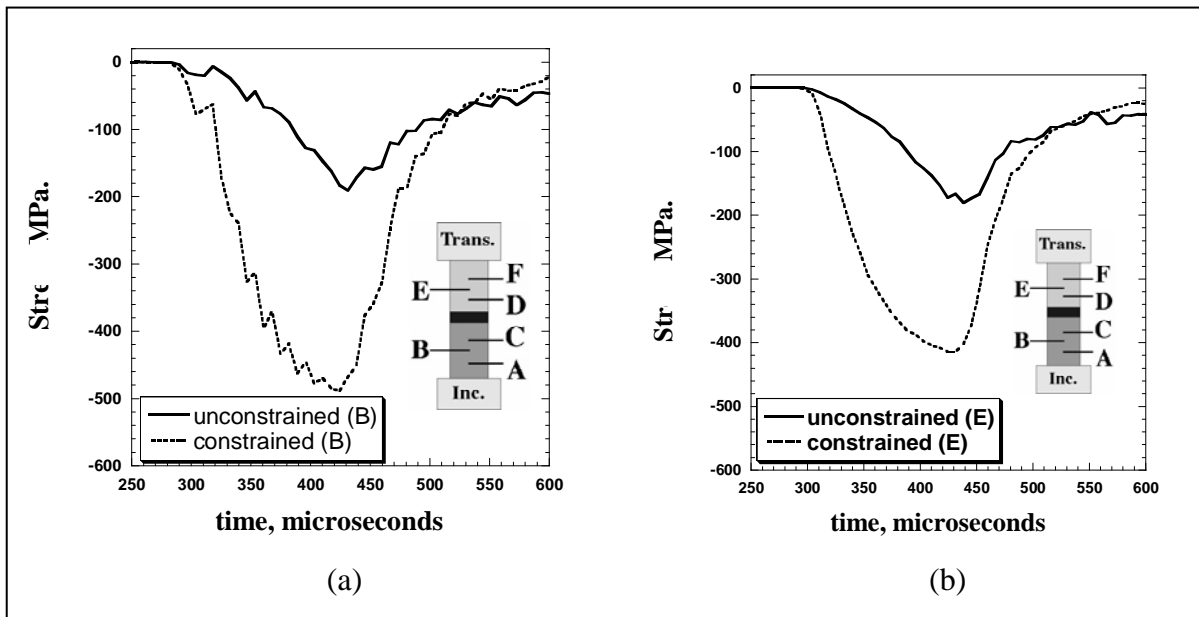


Figure 9. Calculated data from ceramic/rubber/composite ($V = 16$ m/s). Stress measured on (a) ceramic and (b) composite.

stress measured from two strain gages close to incident bar and rubber interfaces of the sample for the ceramic and, similarly, for the composite. (During this particular experiment, the gage close to the transmitter bar interface on the composite broke off due to the high strain at this location. As a result, only the initial portion of the stress read-out can be recorded for this gage.) Figure 11b shows the z-stress in the ceramic layer, calculated at two elements, at almost the same

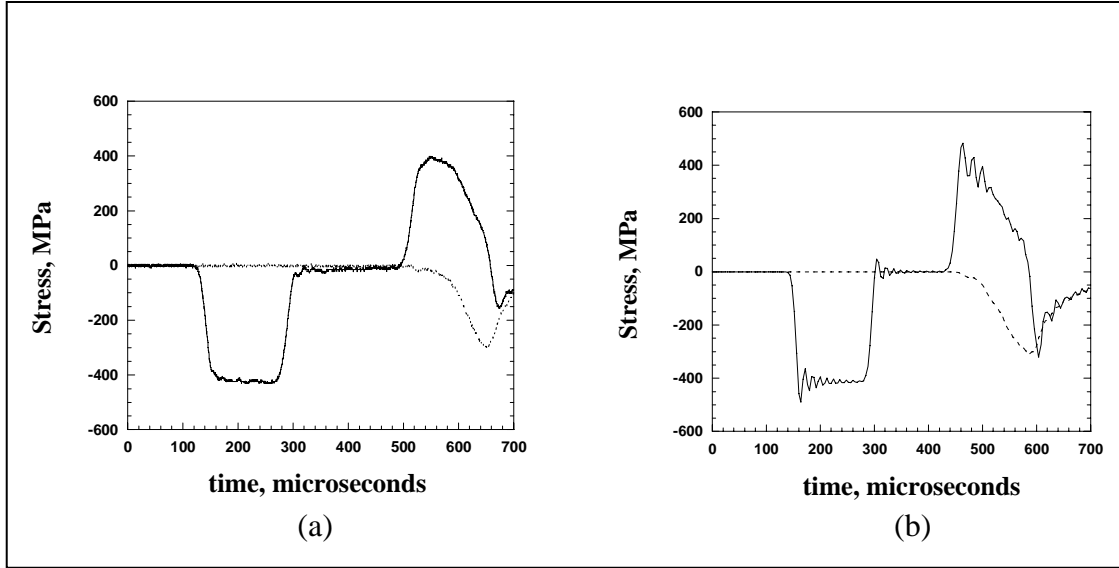


Figure 10. Stress on the incident and transmitter bars during a test at 20.5 m/s on an unconstrained ceramic/rubber/composite: (a) experimental and (b) calculated.

location as the experimental case. It can be clearly seen that during the course of testing the material experiences a nonuniform stress distribution, with the minimum occurring close to the rubber interlayer. Figure 11d shows calculated data from a similar element on the composite layer. Despite possible limitations of the present model due to damage initiation, it is seen that the calculated maximum stress levels are still very close to the measured stresses.

Figure 12a and b shows measured and calculated data from the incident and transmitter bars for the constrained rubber case. For these samples, agreement between the experimental and numerical data is again less close than for the unconstrained case at lower striker bar velocities (figures 6b and 8b) and the peak value of the reflected pulse has again been somewhat overestimated.

Figure 13 presents experimental and numerical data from the individual ceramic and composite layers. Figure 13a and c shows the actual stress measured from a single strain gage at midlength of the sample for the ceramic and composite. Similarly, figure 13b and d shows numerical data from the individual ceramic and composite layers at comparable locations. Comparison with figure 11 shows that a major effect of constraint is to broaden the principal peaks, resulting in the components remaining longer at these higher stress levels. The maximum stress is still experienced in each case ($\sim 140 \mu\text{s}$) after the initial impact. Even for this high-impact velocity, the model still captures the general features of wave propagation and the form of stress distribution. The agreement in terms of absolute values of stress can also be clearly seen from the figures.

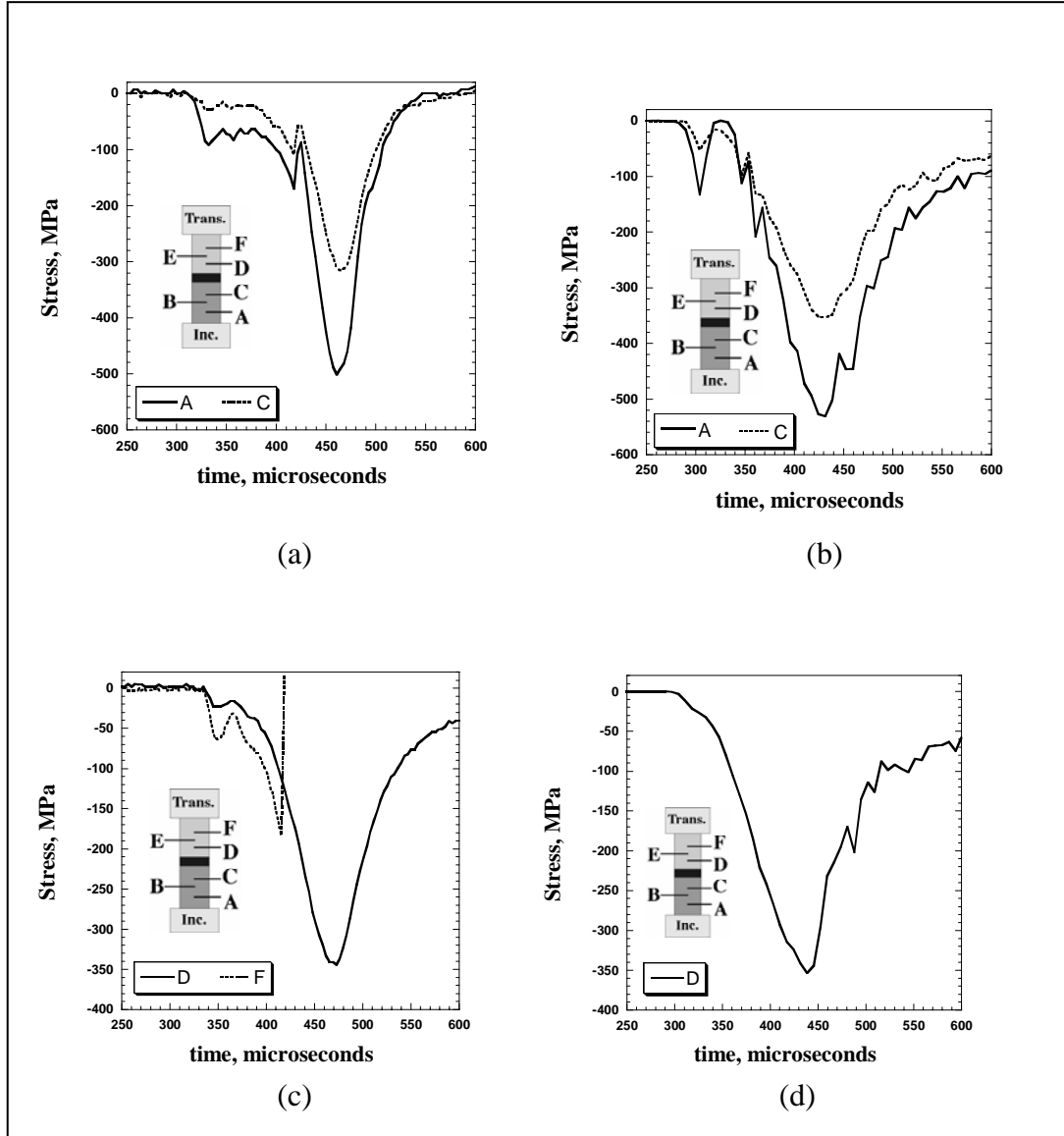


Figure 11. Stress measured on ceramic: (a) experimental and (b) calculated. Stress measured on composite: (c) experimental and (d) calculated (unconstrained-rubber, $V = 20.5$ m/s).

4. Discussion

The main objective of the present work was to investigate experimentally and numerically the effects of lateral confinement of the rubber interlayer on the wave propagation characteristics of the multilayer material over a range of impact velocities. The SHPB is a convenient tool for high strain rate testing of homogeneous elastic/plastic materials, but direct interpretation of SHPB data is not possible for materials which are nonlinear or of very low or very high

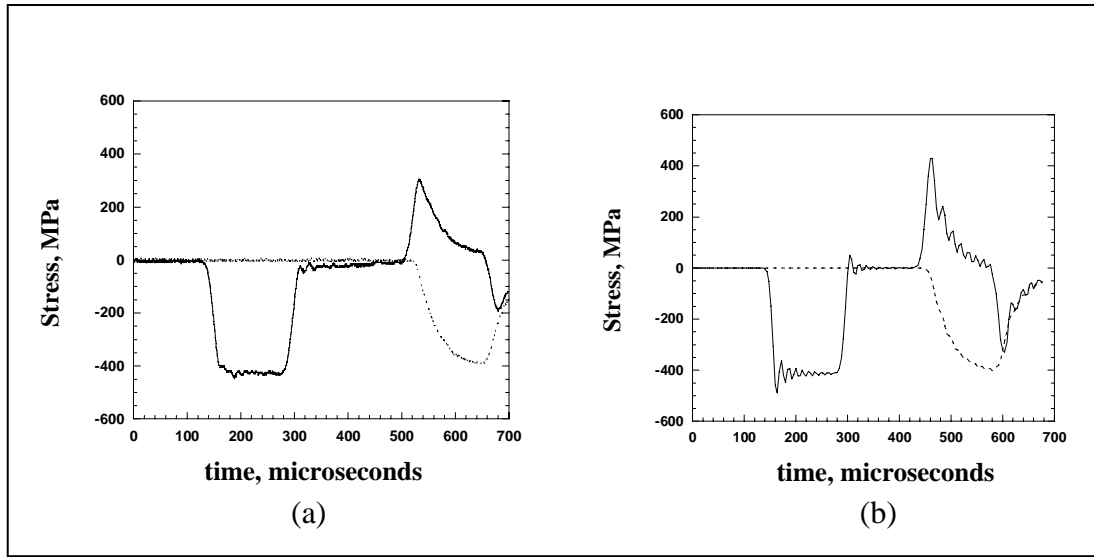


Figure 12. Stress measured on the incident and transmitter bars during a test at 20.5 m/s on a constrained ceramic/rubber/composite: (a) experimental and (b) calculated.

impedance relative to the bars, or anisotropic, or composed of several layers of distinctly different materials as in the case studied here. However, if numerical simulation procedures can be developed which satisfactorily reproduce the output data of SHPB tests, then (a) the tests themselves can be better interpreted and (b) simulations can be carried out with increased confidence.

A previous study (25) showed that there was excellent agreement between numerical data and actual data measured from the incident and transmitter bars for a two-layer ceramic/composite test. That model satisfactorily captured the details of wave transmission. The present three-layer model also satisfactorily captures the details of wave transmission, the general features of wave propagation, stress magnitudes, and the form of stress distribution, and consequently offers considerably enhanced insight into the processes leading to damage generation in such multilayer materials.

The experiments are subject to some limitation because the strain gages average the data over their active gage length, which is typically 0.79 mm as compared with ~0.4 mm for the element size in the model. Also, the measured stress is seen to be very strongly dependent upon the exact placement of the gage within the specimen length. Finally, some of the strain gages mounted on the components could not record all the stress wave history either because of their failure or because of the high strain levels generated in the components (see figure 11c).

Keeping these limitations in mind, it can be appreciated that there is nonetheless good agreement between experimental and numerical data even for the highest impact velocity tests. For example, figure 3a shows experimental data from two gages (5 mm apart) on the ceramic sample

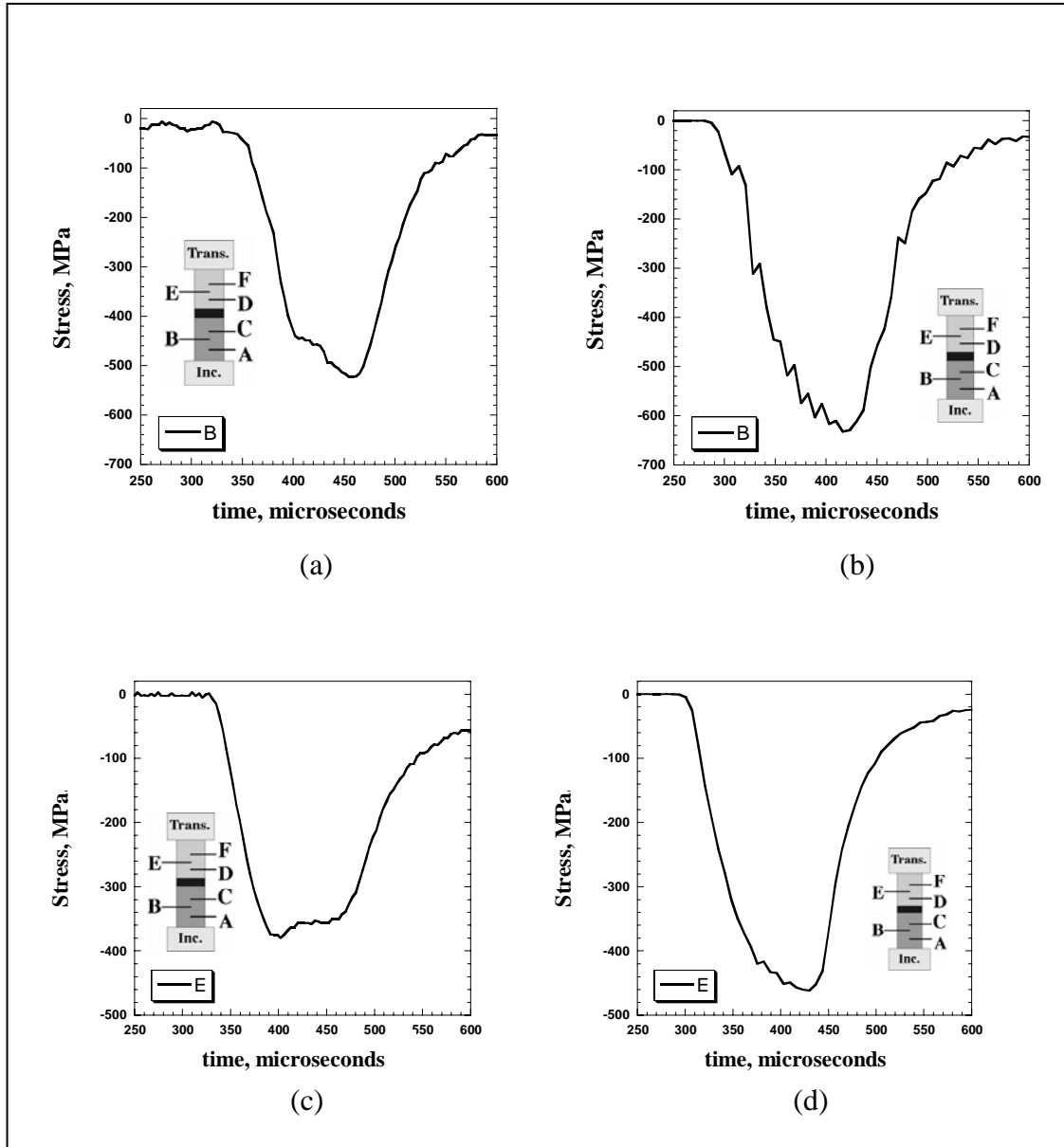


Figure 13. Stress measured on ceramic: (a) experimental and (b) calculated: Stress measured on composite: (c) experimental and (d) calculated (constrained-rubber, $V = 20.5$ m/s).

surface. After the initial stress peak, a second major peak is observed $\sim 150 \mu\text{s}$ later; the numerical data (figure 3b) likewise show the initial and second peak stresses at approximately similar intervals. Furthermore, interpolating to comparable elemental positions, the relative magnitudes are quite similar. The absolute magnitudes of the maximum measured stresses are, however, slightly different, and this is partially due to reasons of data source location and the “averaging” effect of the strain gage size previously mentioned.

Throughout all the current sets of experiments and simulations, close agreement was achieved between the model and experiment. Figures 6a and b and 8a and b show experimental and

calculated data from the Hopkinson bars. Again, the data match closely, showing that LS-DYNA accurately captures the details of wave propagation.

A major conclusion from the present results is that constraint of the rubber layer drastically alters the response of the material, and this is most easily demonstrated by the change of the reflected and transmitted wave shapes. For example, in the case of the lowest velocity tests, instead of reaching a maximum transmitted stress of ~ 60 MPa, as for the unconstrained case, the maximum stress is ~ 210 MPa when constrained (compare figures 2a and 4a). Also the stress rises relatively faster and more uniformly for the second case: basically, lateral confinement of the rubber interlayer increases the wave transmission efficiency between the components.

It can also be seen that the shapes of the wave traveling through the individual ceramic and composite layers are widely different in the unconstrained material, an observation that is confirmed both by experimental measurements and numerical analysis. For example, figure 3 shows that an almost instantaneous and rapid stress increase occurs for the ceramic, followed by further major oscillations, while the composite shows a largely monotonic and gradual increase in stress level. By contrast, in constrained samples, the wave shapes are almost the same for the ceramic and composite (see figure 5) and come to resemble the shape of the wave traveling in the composite. In other words and in common with many other kinetic processes, the component of lowest impedance dominates the process.

Rubber leads to a highly inhomogeneous stress distribution within the components when it is constrained. Generally, the part of the sample close to the unconstrained rubber experiences a reduced stress while the remainder may experience a much higher stress level. However, when the rubber is constrained, the differences in stress level within the components are greatly reduced, although the stress is still by no means homogeneous.

For the higher velocity tests, agreement between the experimental and numerical data is currently slightly less close than for the lower velocity cases (see figures 10–13), for reasons associated with damage evolution. During testing at high velocities, the ceramic frequently shattered, and various damage modes were activated in the composite. Even though the material models used in this study do not include failure parameters, the numerical calculations still capture the general features of wave propagation and the form of stress distribution. Obviously, further refinement along these lines will significantly improve agreement between experiment and model.

In this respect, the present work indicates several avenues for characterizing damage evolution in these, or similar materials at high strain rates since the effects of damage are clearly indicated in the output signals. For example, the principal differences between figures 4a and 10a can be ascribed to the onset of significant damage that alters the (here assumed elastic) properties of the materials. Therefore, the point at which experimental and numerical data begin to diverge probably defines the point at which significant damage begins. It is thus possible to study, by comparison of experimental and numerical data, damage evolution at high strain rate indirectly as a function of strain and strain rate. When coupled with microscopic examination of recovered

material, the present type of data would elucidate the processes and sequence of deformation and fracture events.

This is potentially of great utility because failure criteria are included in recent material models, but a definition of stress and strain levels associated with these events is presently rather imprecise. So, if stress or strain levels associated with the onset of various damage mechanisms can be determined, these could be inserted into the numerical models and provide improved accuracy.

5. Conclusions

The present work has demonstrated the feasibility of modeling stress wave propagation in complex multilayer materials. It has been shown that the effects of confinement of normally low modulus materials can significantly affect their response to wave propagation. Severe stress inhomogeneities and discontinuities may exist in multilayer materials, and these may have serious consequences for the mechanical and other properties. Numerical modeling clearly shows that during Hopkinson bar testing of multilayer materials, stress is not distributed uniformly inside the specimen. The one-dimensional stress state usually assumed for conventional SHPB testing is questionable, and for a complete understanding of the wave propagation, both numerical and experimental results have to be coupled. In this study, both methods were used, and the stress states inside the components were presented. Accuracy will be increased, especially for the high pressure levels, by implementing damage parameters in future material models.

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